

Reduced terrestrial evaporation increases atmospheric water vapor by generating cloud feedbacks

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Abstract

Reduced terrestrial evaporation directly warms the surface by reducing latent cooling, but also indirectly modifies surface climate by altering atmospheric processes. We use a global climate model to explore two end cases of terrestrial evaporation, comparing the climate of SwampLand, a world where land is always fully saturated with water, to that of DesertLand, where land is always completely lacking in soil moisture. When we suppress evaporation to create a desert-like planet, we find that temperatures increase and precipitation decreases in the global mean. We find an increase in atmospheric water vapor over both land and ocean in the DesertLand simulation. Suppressing evaporative cooling over the continents reduces continental cloud cover, allowing more energy input to the surface and increasing surface moist static energy over land. The residence time of atmospheric water vapor increases by about 50 percent. Atmospheric feedbacks such as changes in air temperatures and cloud cover contribute larger changes to the terrestrial surface energy budget than the direct effect of suppressed evaporation alone. Without the cloud feedback, the land surface still warms with suppressed land evaporation, but total atmospheric water vapor decreases, and the anomalous atmospheric circulations over the continents are much shallower than in simulations with cloud changes; that is, the cloud feedback changes the sign of the water vapor response. This highlights the importance of accounting for atmospheric feedbacks when exploring land surface change impacts on the climate system.

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1 Introduction

Changes in terrestrial evapotranspiration directly impact climate. Reducing evaporation from the land surface has a direct warming effect by reducing energy loss from the land surface (Shukla & Mintz, 1982; Fraedrich et al., 1999; Davin et al., 2010; Laguë et al., 2019). However, in idealized continental configurations with large land masses, reducing terrestrial evaporation can instead drive terrestrial cooling by reducing atmospheric water vapor concentrations and the strength of the water vapor greenhouse effect (Laguë, Pietschnig, et al., 2021). Changes in atmospheric temperatures and water vapor driven by changes in terrestrial evaporation are not only important for terrestrial surface climate, but also for the global atmospheric energy budget, as changes in atmospheric temperatures, moisture content, and cloud cover driven by terrestrial processes alter the global radiative balance of the atmosphere (Swann et al., 2010, 2012; Boos & Korty, 2016; Laguë & Swann, 2016; Laguë, Swann, & Boos, 2021).

Terrestrial processes are an integral part of the global water cycle. Water evaporates from the oceans, is transported by the atmosphere, and falls as precipitation over the land. Water on the land surface is evaporated or transpired to the atmosphere, stored as ground water, or returned to the ocean as runoff. Terrestrial evapotranspiration is determined by a combination of soil moisture, atmospheric demand for water, terrestrial re-distribution of water and the evaporative properties of vegetation and soils (Monteith, 1965; Bonan, 2008; Eltahir & Bras, 1996), while large-scale climate features, topography, and soil properties modulate soil water available for evapotranspiration (see Kottek et al., 2006, and references therein).

Terrestrial evaporation can be limited both by the availability of water to evaporate and by energy (Budyko, 1961; Vargas Zeppetello et al., 2019). Vegetation can directly modulate transpiration by opening and closing stomata (Jones, 1998; Sellers et al., 1996; Pielke et al., 1998); transpiration can change with both water availability and vegetation controls on leaf area (Bonan, 2008), stomatal properties (Ball et al., 1987; Medlyn et al., 2011), and root depth (Lai & Katul, 2000). Evapotranspiration (ET) has changed over the historical period (Hobeichi et al., 2021), with some regions (e.g. tropical Africa) exhibiting a negative trend in ET from 1980-2018, while other regions (e.g. Europe) show a positive trend. Earth system models project more changes in the future (Swann et al., 2016; Berg et al., 2016), though the response is complex and varies across models, with stomatal closure in response to increased atmospheric CO₂ acting to reduce ET and increased leaf area and atmospheric demand for water acting to increase ET. Changes in evapotranspiration are an expected result of land use change (Wang et al., 2021), vegetation responses to increased atmospheric CO₂ (Field et al., 1995; Sellers et al., 1996; Norby et al., 2010; Donohue et al., 2013; Lemordant et al., 2018), and climate change (Swann et al., 2016; Berg et al., 2016; Collins et al., 2013).

How changes in terrestrial evaporation relate to the water cycle both regionally and globally remains an area of active research (Swann et al., 2016; Berg et al., 2016; Koster et al., 2006; Dirmeyer, 2011, 1994, 2006; Byrne & O’Gorman, 2015, 2016; Seneviratne et al., 2010; Laguë, Pietschnig, et al., 2021). Total atmospheric water vapor is projected to increase in simulations of global warming (Sherwood et al., 2010); while relative humidity over the oceans is expected to remain roughly constant, relative humidity over land is expected to decrease (O’Gorman & Muller, 2010a; Byrne & O’Gorman, 2016). Independent of the radiative effects of CO₂, plant responses to increased atmospheric CO₂ are projected to reduce near-surface relative humidity on land (Swann et al., 2016).

In this study, we explore the effect of extreme end-cases of terrestrial evaporation on global climate for the modern continental configuration: land that is always fully saturated (all land looks like a swamp), versus land that is always fully desiccated (all land looks like a desert). We find that fully suppressing terrestrial evaporation leads to increased water vapor concentrations throughout the atmospheric column over most con-

81 tinent and ocean regions. While terrestrial relative humidity decreases with suppressed
 82 evaporation, strong cloud feedbacks enhance the energy content and specific humidity
 83 of air over land.

84 2 Methods

85 We conduct experiments using two climate models to study how changes in land
 86 evaporation impact the atmosphere. We use a radiative kernel to decompose the impact
 87 of temperature, moisture, cloud, and albedo changes on the atmospheric energy budget.

88 2.1 Models

89 We use a modified version of the Community Earth System Model (CESM) (Hur-
 90 rell et al., 2013), comprised of the Community Atmosphere Model v. 5 (CAM5) coupled
 91 to a slab ocean model (Neale et al., 2012), the CICE5 interactive sea ice model (Bailey
 92 et al., 2018), and the Simple Land Interface Model (SLIM) (Laguë et al., 2019). Sim-
 93 ulations are run at 2.5° resolution for 50 years, with the first 20 years discarded to al-
 94 low for model spin-up. After 20 years, there is < 0.1 K drift in global mean surface tem-
 95 peratures (Fig. 1); the top of atmosphere energy imbalance in these near-equilibrium sim-
 96 ulations is near-zero (≈ 0.3 W/m²).

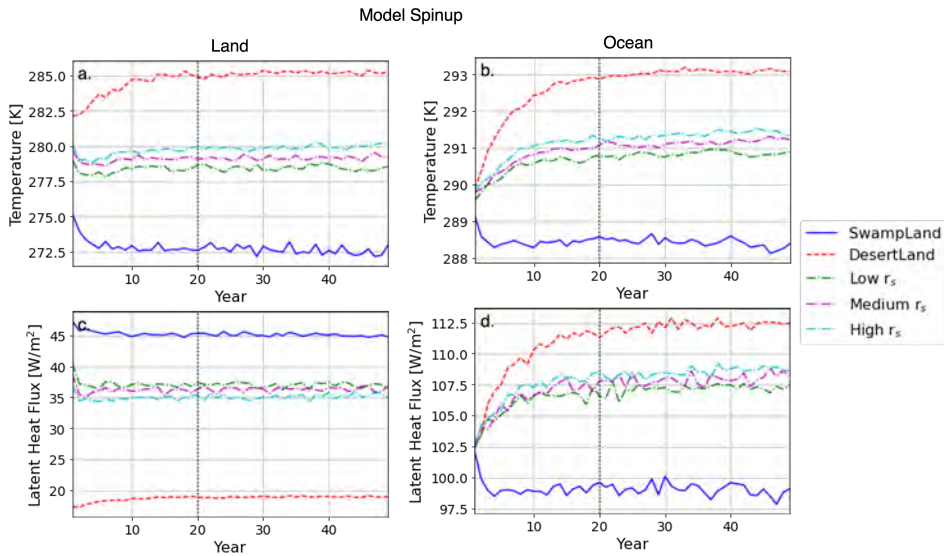


Figure 1. Annual mean, spatially averaged surface temperature (top) and latent heat flux (bottom) over land (left) and ocean (right) regions for the five CESM simulations explored in this study. Simulations have reached equilibrium prior to year 20.

97 The slab ocean has prescribed, seasonally and spatially varying pre-industrial ocean
 98 heat transport and heat capacity (slab or mixed-layer “depth”) which is identical between
 99 simulations and from year-to-year within the same simulation. This allows sea surface
 100 temperatures to evolve in response to forcings, but does not allow for changes in ocean
 101 heat transport. The mixed-layer depth and heat transport values used in these simula-
 102 tions are calculated from the dynamic ocean component of the pre-industrial control sim-
 103 ulation of the fully coupled CESM 1.2 model (as used in Garcia et al., 2016; Swann et
 104 al., 2018; Laguë et al., 2019). We use slab ocean values from a pre-industrial vs. present-
 105 day simulation because the pre-industrial climate is in equilibrium, while the ocean acts

106 as a net energy sink in the present day climate. The ocean is a large source of variability
107 in the Earth system, and the use of a slab ocean model allows us to focus on the at-
108 mospheric response to land surface changes, as the internal variability of the slab ocean
109 model is small, allowing us to use a single simulation for each experiment rather than
110 an ensemble. To explicitly demonstrate this, we run the initial 10 years of the two sim-
111 ulations described below (“SwampLand” and “DesertLand”) with three ensemble mem-
112 bers, each with initial land surface temperatures perturbed by $1e-6$ K; the spread be-
113 tween ensemble members is very small compared to the difference between the two ex-
114 periments (Fig. S1). The default ocean albedo parameterization is used, where ocean albedo
115 varies with solar declination, with default values of 0.06 for direct and 0.07 for diffuse
116 radiation.

117 SLIM is used to allow us to directly control the physical properties of the land sur-
118 face in a way that is difficult with complex land surface models. Hydrology is represented
119 using a bucket model, with the resistance to evaporation calculated as a function of how
120 much water is in the bucket as well as an additional user-prescribed resistance. A sim-
121 ple snow model allows for snow-albedo feedbacks on the land surface. The idealized land
122 surface model allows us to artificially control terrestrial water availability without alter-
123 ing other aspects of the land surface. In contrast, if we were to override soil moisture in
124 a complex land surface model like CLM5 (Lawrence et al., 2019), this would have follow-
125 on impacts on leaf area, plant hydraulics, and the carbon cycle, which would in turn have
126 follow-on impacts on albedo and surface evaporative resistance. In order to isolate the
127 effect of surface water availability on climate, we need to leverage an idealized land sur-
128 face model like SLIM.

129 We also want to test the response of the climate system to changes in land evap-
130 oration without cloud responses, which are a large source of model uncertainty. We could
131 force clouds in CESM to be transparent to radiation, but this would lead to an unrea-
132 sonably dark top of atmosphere albedo (unless we *also* modified surface albedo) and a
133 much hotter base-state climate than CESM produces with the normal cloud parameter-
134 ization. Instead, we conduct two additional simulations using Isca (Vallis et al., 2018),
135 an idealized global circulation model which has radiatively interactive water vapor and
136 produces precipitation, but where clouds are “invisible” to radiation in the standard, tested
137 configuration. Simulations use a T42 grid, a slab ocean with a 20m mixed layer depth
138 and prescribed modern heat transport, a bucket model for land hydrology with a heat
139 capacity equal to that of a 2m ocean mixed layer, no snow albedo feedbacks, the mod-
140 ern continental configuration, realistic topography, and the RRTM radiation scheme (Clough
141 et al., 2005). The albedo of the surface is set to be much higher than realistic values (0.25
142 for ocean and 0.325 for land), to generate a reasonable climate and top of atmosphere
143 albedo (Thomson & Vallis, 2019; Geen et al., 2018).

144 2.2 Simulations

145 We consider five CESM simulations in this study, primarily focusing on two end-
146 members of wet/dry land. Simulations have globally uniform land surface properties: all
147 non-glaciated points on the land surface have the same albedo, aerodynamic roughness,
148 capacity to hold water, and evaporative resistance. Glaciated gridcells (Greenland and
149 Antarctica) have surface properties associated with ice, and do not vary between sim-
150 ulations (see Laguë et al., 2019, for details). The snow-free albedo of the land surface
151 is set to 0.2 in visible wavelengths and 0.3 in the near-infrared, while the aerodynamic
152 roughness is set to 0.1 m. While the base-state climate differs regionally in these sim-
153 ulations with globally uniform land surface properties compared to simulations with present-
154 day land surface properties (see Laguë et al., 2019), we choose to use uniform, idealized
155 land surface properties in order to isolate the sensitivity of the atmosphere to only the
156 change in water availability at the land surface, without introducing the added dimen-

157 sion of imposing the difference in water availability onto different land surface types (e.g.
158 land with different albedo).

159 In DesertLand, the first of our extreme simulations, the capacity of the land to hold
160 water is reduced to 20 mm everywhere (compared to a typical value of ≈ 200 mm), and
161 the resistance to evaporation is set to 100,000 s/m (compared to a typical value of ≈ 100
162 s/m); this effectively turns off evaporation from the land surface, regardless of precip-
163 itation or the atmospheric demand for water. DesertLand can physically be thought of
164 as land free of vegetation with extremely well draining soils, such that all precipitation
165 that falls on the land is quickly transferred into below ground aquifers or sub-surface runoff
166 and returned to the oceans. In SwampLand, the second extreme simulation, the land sur-
167 face is forced to be fully saturated with water at every time step. Land always has 200
168 mm of water available for evaporation, regardless of the precipitation or evaporation rates
169 at each point. SwampLand can be thought of as land with a high water table and un-
170 limited ground water supply. Physically, this is comparable to swampy regions on the
171 modern land surface, but in this idealized simulation, these swamps are imposed over
172 the entire non-glaciated land surface, regardless of elevation, slope, or distance from a
173 water body.

174 Three additional simulations with interactive surface hydrology are also briefly con-
175 sidered, differing in their prescribed evaporative resistance: 30 s/m (“low” for low re-
176 sistance), 100 s/m (“medium” for medium resistance), and 200 s/m (“high” for high re-
177 sistance). Each simulation has the capacity to hold 200 mm of water at each non-glaciated
178 land surface point, but the amount of water actually on the land surface is interactively
179 simulated by the model based on precipitation and evaporation at each location.

180 Two additional simulations, similar to DesertLand and SwampLand, are conducted
181 in the Isca model. In the DesertLand Isca simulation, the bucket capacity is set to 0.01
182 mm, while the bucket capacity in SwampLand is set to 150 mm and is “topped up” to
183 always have 150 mm of water available at all land points at each time step.

184 **2.3 Radiative Kernel**

185 Changing surface evaporation between simulations alters atmospheric and surface
186 temperatures, water vapor, cloud cover, and snow and ice extent. To isolate the individ-
187 ual contributions of each of these responses to the atmospheric energy budget, we use
188 a radiative kernel for the CAM5 atmospheric model (Pendergrass et al., 2018) and fol-
189 low the procedure introduced in Laguë, Swann, & Boos (2021), which we summarize here.

190 The radiative kernel provides the change in top-of-atmosphere (TOA) net short-
191 wave and longwave radiation under both “full-sky” (including the effects of clouds) and
192 “clear-sky” (without the effects of clouds) conditions that result from independent changes
193 in the following quantities: surface albedo; surface temperature; air temperature at each
194 level of the atmosphere; and specific humidity given a unit change in air temperature at
195 constant relative humidity at each level of the atmosphere. The kernel also provides the
196 change in downwelling shortwave and longwave radiation at the surface associated with
197 each of these perturbations. Multiplying our simulated change in temperature, water va-
198 por, etc. by the kernel returns the effect of that change on TOA or surface radiative fluxes.

199 We mask out differences in the stratosphere between simulations (as in Pendergrass
200 et al., 2018; Shell et al., 2008), and multiply the water vapor kernel by the change in the
201 natural logarithm of the simulated change in water vapor. We apply the clear-sky lin-
202 earity test (Vial et al., 2013) to our most extreme simulations (SwampLand and Desert-
203 Land) and find generally excellent agreement between the change in TOA fluxes simu-
204 lated by the full model and those predicted by the radiative kernel (Fig. 2). The excep-
205 tion is in the deep tropics, where the kernel and model are qualitatively similar, but dis-
206 agree by a few W/m^2 ; these changes are primarily driven by disagreements in the trop-

207 ical Atlantic and over tropical Africa (not shown). Note that all simulations here have
 208 a net input of energy into the TOA in the low latitudes and a net removal of energy from
 209 the TOA in the high latitudes under both clear-sky and full-sky conditions; Fig. 2 shows
 210 the *difference* in the TOA energy balance between two simulations. In the high latitudes,
 211 the positive values on this graph show that DesertLand is less negative than SwampLand,
 212 i.e. DesertLand is losing less energy at this latitude than SwampLand is. The negative
 213 values in the low latitudes show that DesertLand is less positive than SwampLand, i.e.
 214 that SwampLand is absorbing more total net energy into the Earth System at low lat-
 215 itudes compared to DesertLand. When we show results of calculations using the radiative
 216 kernel, any non-linearities or residuals in the radiative kernel are necessarily included
 217 in the cloud term; however, the excellent agreement in the clear-sky linearity test gives
 218 us confidence that such residuals are small.

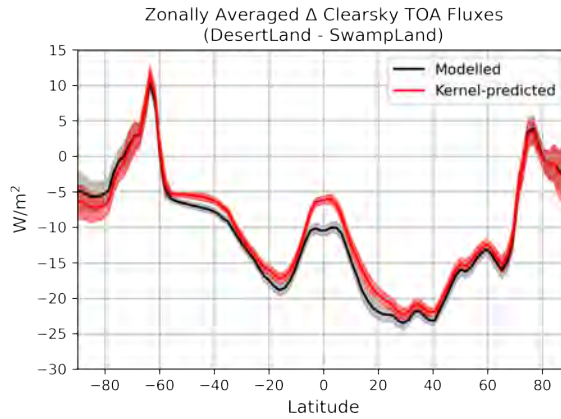


Figure 2. Zonal mean change (DesertLand - SwampLand) in TOA clear-sky radiation directly from the model (gray) and predicted by the clear-sky radiative kernel (red). Shading indicates $\pm 1\sigma$ of inter-annual standard deviation.

219 Analysis was conducted with the Python programming language, primarily with
 220 the NumPy (Harris et al., 2020) and xarray (Hoyer & Hamman, 2017) packages, using
 221 the JupyterHub (<https://jupyter.org/>) service on the Cheyenne computing system
 222 (Computational and Information Systems Laboratory, 2019). When statistical signifi-
 223 cance is shown on maps and vertical cross-sections, a value is deemed statistically sig-
 224 nificant if the p-values calculated from a Student’s t-test pass a false discovery rate of
 225 0.15 (following Wilks, 2016). Uncertainty intervals indicate $\pm 1\sigma$ of inter-annual stan-
 226 dard deviation.

227 3 Results & Discussion

228 3.1 Column water vapor increases with suppressed land evaporation

229 As expected, SwampLand—the simulation with perpetually saturated land—has
 230 the largest terrestrial evaporation of the simulations considered, and the lowest average
 231 land surface temperatures (Fig. 3). In contrast, DesertLand—the simulation with per-
 232 petually suppressed land evaporation—has the lowest terrestrial evaporation (by design),
 233 and the warmest surface temperatures both on land and globally (Fig. 3a,b). The three
 234 simulations with intermediate values of terrestrial evaporative resistance lay between Swamp-
 235 Land and DesertLand in terms of evaporation, surface temperatures, and total atmo-
 236 spheric water vapor. There is a strong linear relationship across the simulations between

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terrestrial evaporation and terrestrial surface temperature, and between global mean surface temperatures and total atmospheric water vapor (Fig. 3c,d).

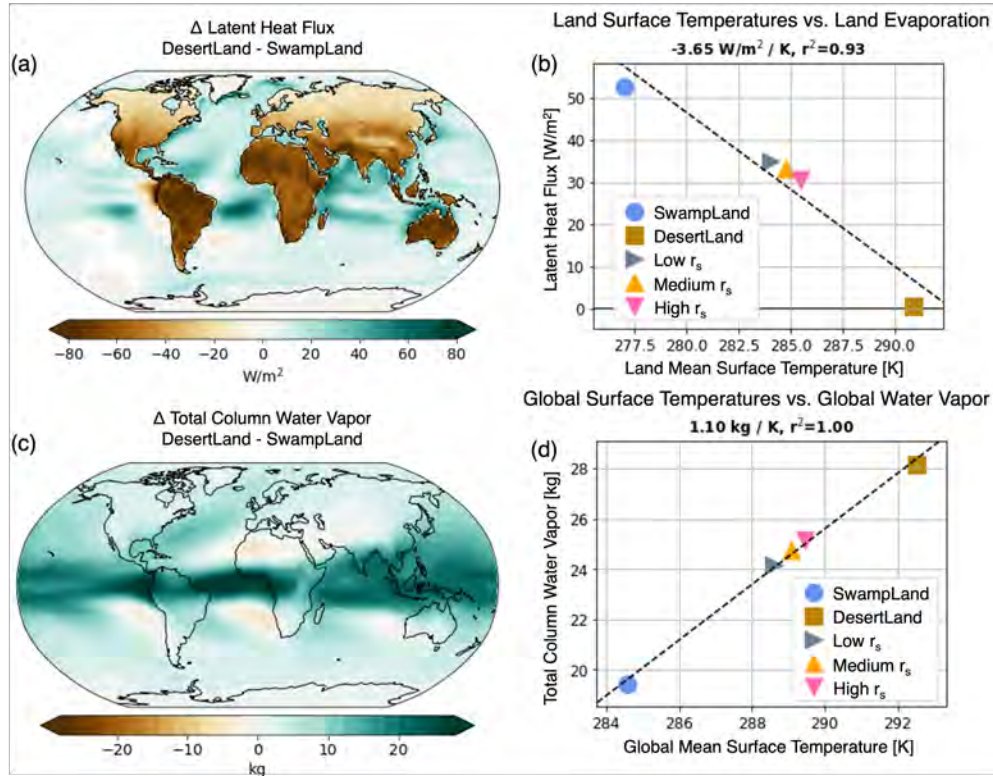


Figure 3. Annual mean change in latent heat flux (a) and total column water vapor (c) for the DesertLand - SwampLand simulations. Scatter plots showing the relationship between annual mean (b) land surface temperature and terrestrial evaporation and (d) global mean surface temperature and global mean total column water vapor. In a/c, only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.

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DesertLand has the most atmospheric water vapor, despite having suppressed land evaporation (Figs. 3). The planet as a whole is not water limited in the modern continental configuration, so ocean evaporation increases in the DesertLand simulation (Fig. 4a). However, this only partially compensates for the reduction in land evaporation, so there is less surface evaporation in the global mean in DesertLand. Precipitation over both land and ocean is accordingly reduced in DesertLand, but the atmosphere has more total water vapor in DesertLand than SwampLand (Fig. 4c). This is true of both land and ocean regions in the lower and upper troposphere, except for a drying of some subtropical regions and in the lower troposphere of inland continental regions (Figs. 5, 6).

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The reduction in global mean precipitation and increase in global mean water vapor content together imply an increase in the residence time of atmospheric water vapor. Specifically, this residence time has been defined as the ratio of global mean precipitable water Q to global mean precipitation P (Trenberth, 1998),

$$\tau \equiv \frac{Q}{P} \quad (1)$$

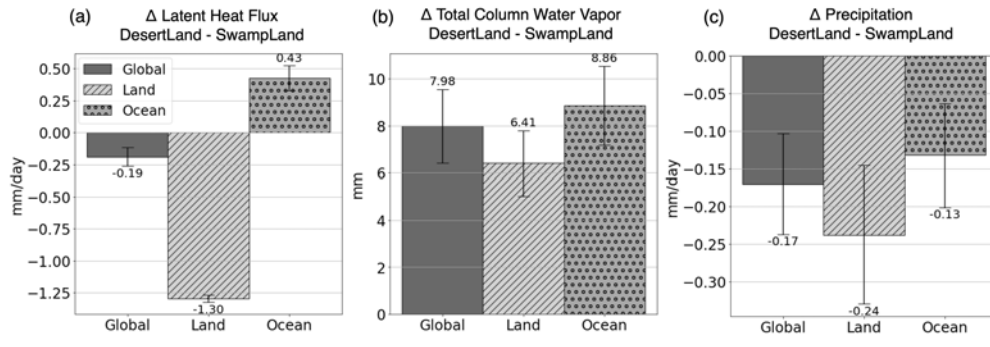


Figure 4. Annual mean change (DesertLand - SwampLand) in area-weighted mean (a) latent heat flux [mm/day], (b) column water vapor [mm], and (c) precipitation [mm/day], separated into the global (solid), land (hatched), and ocean (dotted) contributions. Error bars indicate $\pm 1\sigma$ of inter-annual standard deviation. Note that as there is more ocean than land area, bars within each subplot do not sum directly.

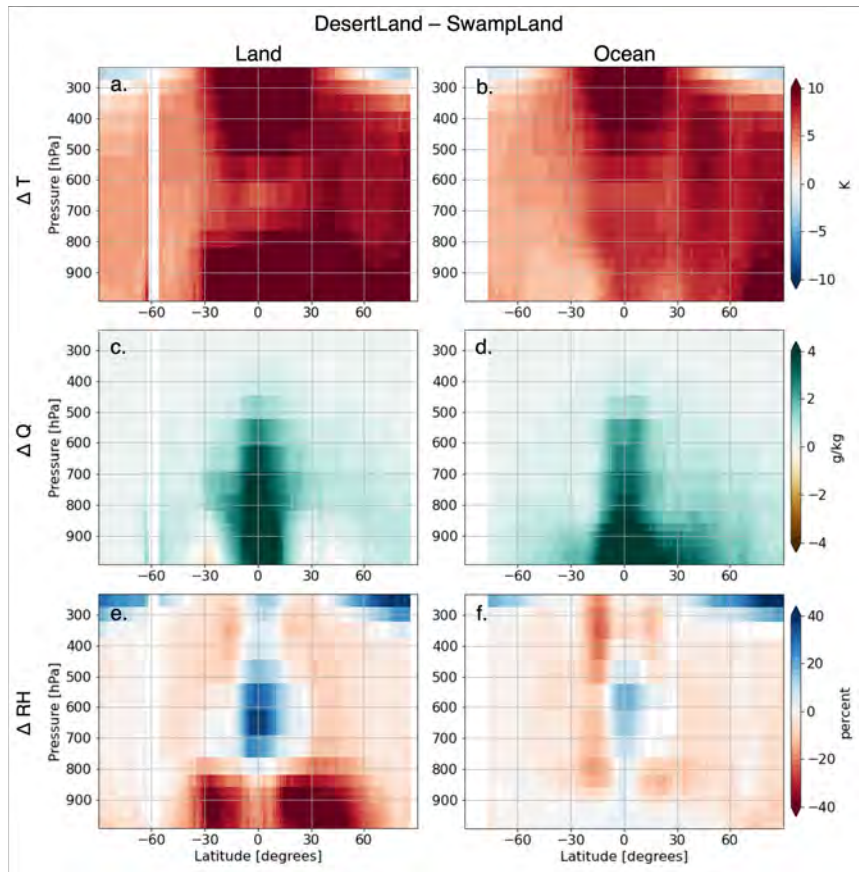


Figure 5. Zonal mean vertical cross sections of the change in temperature (T [K], top), specific humidity (Q [g/kg], middle), and relative humidity (RH [%], bottom) over land regions (left) and ocean regions (right). Only changes that pass a statistical test are shown, where values are significant if the p -values calculated from a student's t -test pass a false discovery rate of 0.15.

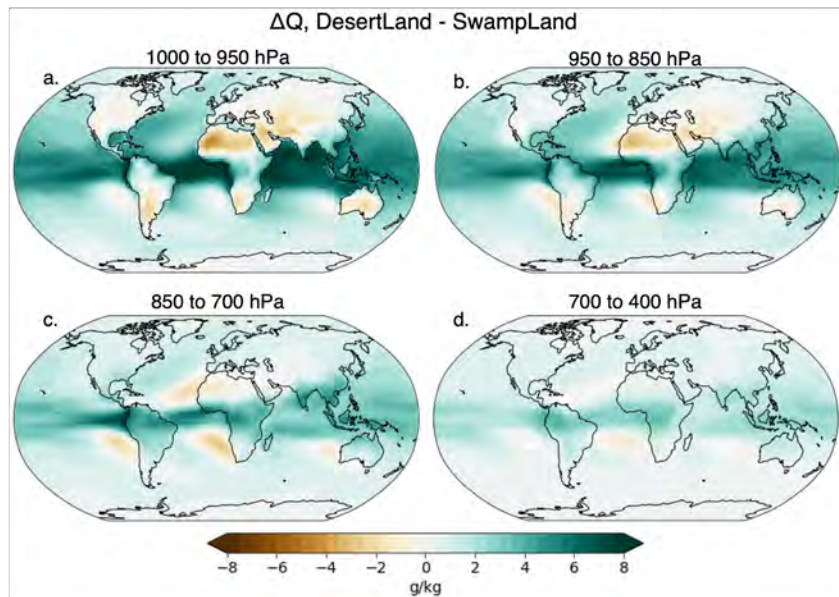


Figure 6. Change (DesertLand - SwampLand) in specific humidity [g/kg] from (a) 1000-950 hPa, (b) 950-850 hPa, (c) 850-700 hPa, and (d) 700-450 hPa. Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.

252 Here we find that τ increases from 6.7 days in SwampLand to 10.2 days in DesertLand.
 253 One may alternatively interpret this change as a reduction in the convective mass flux
 254 that transports water vapor vertically until it condenses and precipitates (Held & So-
 255 den, 2006). Thus, while a reduction in land evaporation is expected to produce a tran-
 256 sient reduction in local atmospheric water vapor, changes in the precipitating atmospheric
 257 circulation dominate to allow more water to accumulate in the atmosphere and then be
 258 maintained at that higher level. Locally, water vapor over land is maintained by a bal-
 259 ance between local evaporation, local precipitation, and the convergence of water by large-
 260 scale winds; the latter two are typically large compared to the former, making possible
 261 indirect effects of a surface evaporative forcing. Furthermore, Sun & Wang (2022) found
 262 a reduction in precipitation intensity in hot weather in moisture limited regions (e.g. over
 263 land) as a result of an increased saturation deficit.

264 3.2 Cloud feedbacks enhance energy input over land

265 We now describe how suppression of surface evaporation produces a reduction in
 266 low cloud cover (Fig. 7), increasing the energy absorbed by land and the precipitating
 267 large-scale circulation driven by that energy source. Reductions in cloud cover driven
 268 by suppressed terrestrial evaporation lead to an additional ≈ 37 W/m² of shortwave ra-
 269 diation absorbed by the surface (Fig. 8). The largest reductions occur in low clouds over
 270 land, and are thus consistent with a local response to the land evaporative forcing. The
 271 reduced cloud cover also makes it easier for longwave radiation emitted by land to exit
 272 the top of the atmosphere, but this is smaller in magnitude than the shortwave cloud
 273 effect (as expected given the low altitude of the cloud changes), yielding a net positive
 274 heating of land by cloud radiative effects of about 11 W/m². Other components of the
 275 surface energy budget are discussed below, when we use the model's radiative kernel to
 276 decompose the net change into different physical components.

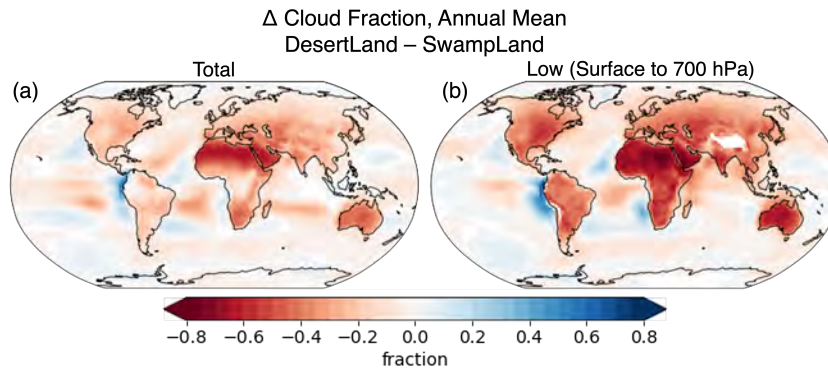


Figure 7. Annual mean change (DesertLand - SwampLand) in total cloud fraction (left) and low cloud fraction (right) for the CESM simulations. Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.

277 Reductions in low cloud cover over land as a result of suppressing terrestrial evap-
 278 oration also impact the TOA energy budget (Fig. 9). In particular, loss of low cloud cover
 279 over land lowers the planetary albedo; this results in more energy absorbed at the TOA
 280 and in turn should increase global mean temperatures. Increased temperatures should,
 281 in turn, lead to an increase in water vapor following the Clausius-Clapeyron relationship.
 282 That is, the reduction in low cloud cover alone should lead to increased atmospheric tem-
 283 peratures and water vapor. Without the reduction in low cloud cover, suppressing ter-
 284 restrial evaporation would not necessarily lead to any increased energy into the Earth
 285 system at the top of the atmosphere, and would lead to weaker warming at the surface
 286 as the only warming mechanism would be the reduction in latent cooling, with no ad-
 287 ditional warming from increased solar radiation due to reduced cloud cover. Indeed, we
 288 see this later when we conduct similar simulations in a model without cloud cover.

289 We roughly estimate the effect of the change in clouds on global mean tempera-
 290 tures and water vapor as follows: in the global mean, the change in clouds leads to a 12.8
 291 W/m^2 increase in energy into the Earth system at the TOA (14.8 W/m^2 from shortwave
 292 radiation, -2 W/m^2 from longwave radiation; Fig. 9). Using a climate response param-
 293 eter of $1.5 \text{ W/m}^2/\text{K}$ (from Gregory (2004)'s estimates for increased TOA insolation), this
 294 would result in a roughly 8.5 K temperature increase, close to the actual global mean
 295 temperature increase in our simulations of roughly 8 K (from a global mean of 283.5 to
 296 291 K; Fig. 3d). Assuming that column water vapor scales with surface air temperature
 297 at a rate of roughly 7%/K per the Clausius-Clapeyron relationship (O'Gorman & Muller,
 298 2010b; Held & Soden, 2006), we would expect water vapor to increase roughly 70% (the
 299 Clausius-Clapeyron equation is exponential, so the relative increase in saturation vapor
 300 pressure for the relatively large warming of 8 K is substantially larger than the infinites-
 301 imal rate of change of 7%/K); the actual increase in total atmospheric water vapor in
 302 our simulations was roughly 40% (from a global mean of 19 to 28 kg; Fig. 3d), with re-
 303 ductions in tropical lower-tropospheric relative humidity over land that peak near -40%
 304 (Fig. 5e). Thus, this simple argument provides similar order of magnitude changes in tem-
 305 peratures and water vapor as our full simulations, but with a nearly factor-of-two over-
 306 estimate in the vapor increase because the radiative forcing (a reduction in low clouds)
 307 is driven by lower-tropospheric drying. We note that there is a lot of uncertainty in the
 308 value of the climate response parameter, both across models and across forcing mech-
 309 anisms (e.g. CO_2 , insolation, etc.) within a single model (Gregory, 2004; Hansen et al.,

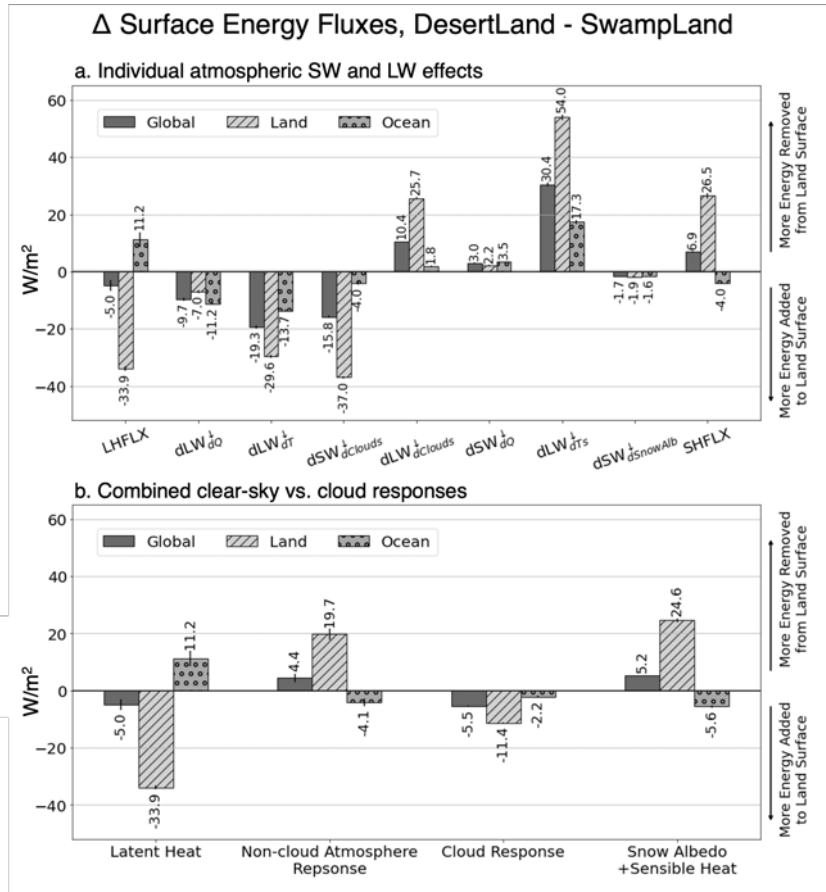


Figure 8. Change in global (solid), land (hatched), and ocean (dotted) annual mean surface fluxes for DesertLand - SwampLand. The breakdown of surface flux changes due to latent (LH-FLX) and sensible (SHFLX) heat, and shortwave (SW) and longwave (LW) radiation are shown. Panel (a) shows the changes in surface shortwave and longwave radiative fluxes decomposed using the radiative kernel into the contributions due to water vapor, atmospheric temperatures, surface temperatures, cloud cover, and surface albedo (i.e. snow changes), as well as the changes due to latent and sensible heat. Panel (b) combines the fluxes into those driven by latent heat flux, the cloud-free atmosphere and Plank response (air temperature, water vapor, and surface temperature), those driven by clouds, and other surface changes (sensible heat flux, albedo). The inter-annual standard deviation is marked by the vertical black lines capping each bar. Downwards (negative) values indicate that the change in the flux leads to more energy into the land surface (i.e. a warming effect), while upwards (positive) values indicate less energy into the land surface or more energy removed from the land surface.

1997), and that 7%/K is not an exact Clausius-Clapeyron scaling (O’Gorman & Muller, 2010b).

3.2.1 Decomposing the energy balance with a radiative kernel

The increase in near-surface MSE is driven by a combination of factors. In the absence of evaporative cooling (i.e. the DesertLand simulation), changes in the atmosphere which increase radiative fluxes into the land surface necessarily lead to warming. Over

316 land, suppressed terrestrial evaporation directly increases the energy that must be re-
 317 moved from the surface as sensible heat or longwave radiation (Fig. 8, spatial patterns
 318 shown in Fig. S2). Reducing evaporation results in excess energy available in the land
 319 surface ($\approx 34 \text{ W/m}^2$ averaged over all land areas). The increase in energy into the land
 320 surface from increase downwelling longwave radiation as a result of suppressed land evap-
 321 oration (29.6 W/m^2 from air temperatures and an additional 7 W/m^2 from increased
 322 water vapor) is of a similar magnitude to the increase in energy into the land surface from
 323 suppressed latent heat flux (33.9 W/m^2).

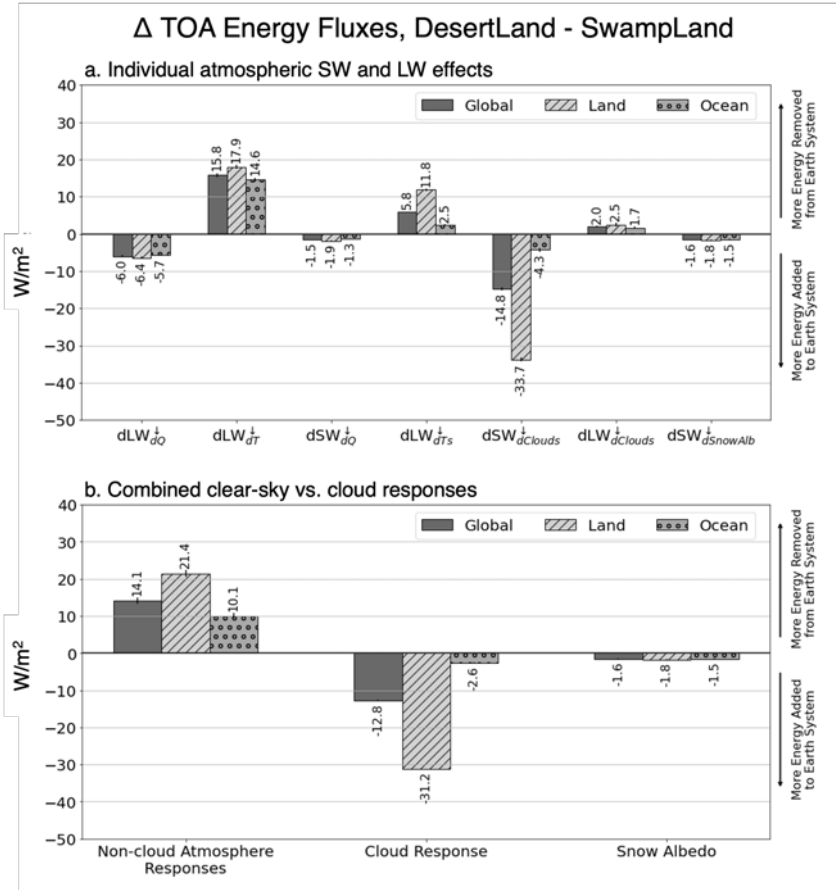


Figure 9. Change in global (solid), land (hatched), and ocean (dotted) annual mean top of atmosphere (TOA) fluxes for DesertLand - SwampLand. The breakdown of toa flux changes due to shortwave (SW) and longwave (LW) radiation are shown. Panel (a) shows the changes in toa shortwave and longwave radiative fluxes decomposed using the radiative kernel into the contributions due to water vapor, atmospheric temperatures, surface temperatures, cloud cover, and surface albedo (i.e. snow changes). Panel (b) combines the fluxes into those driven by the cloud-free atmosphere and Plank response (air temperature, water vapor, and surface temperature), those driven by clouds, and surface albedo. The inter-annual standard deviation is marked by the vertical black lines capping each bar. Downwards (negative) values indicate changes that lead to more energy absorbed by the Earth system at the TOA while upwards (positive) values indicate energy removed from the Earth system at the TOA.

324 The longwave effects of increased water vapor act against the negative longwave
 325 cloud effect (Figs. 8, S2b,d). Increased water vapor itself drives a feedback on the sur-

326 face energy balance, with more water vapor leading to more longwave radiation into the
 327 surface (7.9 W/m^2); this term is of comparable magnitude to (though slightly smaller
 328 than) net cloud radiative effects (11.4 W/m^2) over land; Figs. 8, S2). Previous work has
 329 shown how changes in terrestrial evaporation modulate the water vapor greenhouse ef-
 330 fect; specifically, Laguë, Pietschnig, et al. (2021) show that while reducing land evapo-
 331 ration directly warms the surface, over very large idealized continents, reductions in land
 332 evaporation lead to reduced atmospheric water vapor and drive an overall cooling at the
 333 surface by reducing the water vapor greenhouse effect. In this study (with the modern
 334 Earth’s continental configuration), our results show that land’s control on water vapor
 335 is still an important contribution to the radiation budget, with the changes in surface
 336 and TOA fluxes driven by changes in water vapor of comparable magnitude to the com-
 337 bined shortwave and longwave effects of changes in cloud cover.

338 The combined effect of atmospheric responses to suppressed terrestrial evapora-
 339 tion is a slight increase in longwave energy into the surface, and a larger increase in short-
 340 wave radiation into the surface which, over land areas, is comparable in magnitude to
 341 the increase in surface energy coming from reduced latent cooling. Suppressing terres-
 342 trial evaporation leads to a net radiative flux of roughly 6 W/m^2 out of the land surface
 343 (the sum of all the radiative fluxes in Fig. 8a), which is balanced by the combined changes
 344 in sensible and latent heat flux. While increased sensible heat flux (associated with higher
 345 surface temperatures in DesertLand vs. SwampLand) increases energy removed from the
 346 surface (by 26.5 W/m^2 over land), suppressed latent heat flux over land means 33.9 W/m^2
 347 of energy is not removed from the land surface through evaporation. The response of the
 348 radiative terms of the surface energy budget over the oceans are of the same sign as the
 349 changes over land, though of different magnitude. In contrast, because ocean evapora-
 350 tion increases when terrestrial evaporation is suppressed, this removes energy from the
 351 surface, balanced by reductions in sensible heat flux. Over land areas, the radiative ef-
 352 fects of cloud changes are larger than the combined radiative effects of water vapor and
 353 temperature changes (Fig. 8b).

354 The importance of the cloud feedback is reinforced by considering the change in
 355 top of atmosphere (TOA) energy fluxes. The shortwave effects of reductions in cloud cover
 356 in DesertLand-SwampLand are the single largest contributor to changes in the top of at-
 357 mosphere energy balance over land (Fig. 9, spatial patterns in Fig. S3; the largest cloud
 358 reductions occur in low clouds over land areas (Fig. 7). The difference between the TOA
 359 and surface energy flux anomalies is the anomalous net energy input (NEI) to the at-
 360 mosphere discussed below. Both the DesertLand and SwampLand simulations are in equi-
 361 librium, so the net TOA energy balance (the sum of the bars in Fig. 9) is near zero.

362 **3.2.2 Enhanced energy sources drive tropical ascent**

363 The reduction in low clouds in the Desertland vs. SwampLand simulation lead to
 364 an increase in the total amount of solar energy absorbed at the land surface—energy which
 365 is then emitted back to the atmosphere through the surface energy budget. Specifically,
 366 suppressed terrestrial evaporation leads to an increase over most land regions in the net
 367 energy input to the atmosphere (NEI), which is the sum of radiative and surface turbu-
 368 lent fluxes into the atmosphere through its top and bottom boundaries (Fig. 10d). It also
 369 leads to an increase in near-surface moist static energy (MSE) over most continental re-
 370 gions, especially in the tropics (Fig. 11). We calculate moist static energy as

$$MSE = c_p T + L_v Q + gZ, \quad (2)$$

371 the sum of the dry energy (the heat capacity c_p of dry air multiplied by the air temper-
 372 ature T), potential energy (the gravitational constant of acceleration g multiplied by the
 373 geopotential height Z), and the moist energy (the latent heat of vaporization L_v multi-
 374 plied by water vapor Q) of a static (non-dynamic) parcel of air. Note that while we show
 375 the NEI for clear-sky CESM, the latent and sensible heat fluxes which go into the NEI

376 calculation are from the full CESM model that includes cloud effects on the surface en-
 377 ergy budget; as such, Fig. 10e should be treated with appropriate caution.

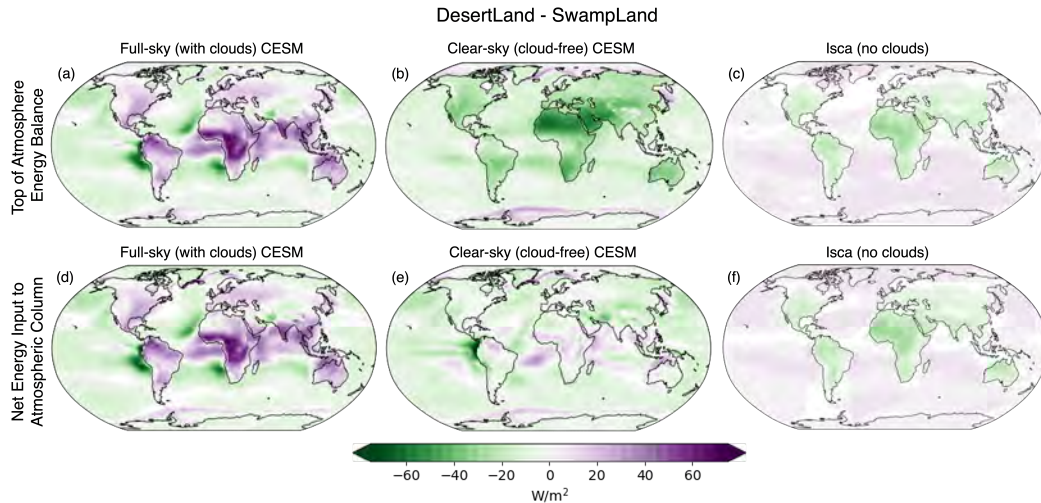


Figure 10. Annual mean change (DesertLand - SwampLand) in the top of atmosphere energy budget (net shortwave - outgoing longwave radiation; top row) and net energy input (NEI: net top of atmosphere - net surface fluxes; bottom row). The full CESM simulation, including cloud radiative effects, is shown in the left column; CESM where radiative fluxes ignore the influence of clouds (“clear-sky” conditions) are shown in the centre column, while Isca, which is always radiatively cloud-free, is shown in the right column. Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student’s t-test pass a false discovery rate of 0.15.

378 Two complementary theoretical frameworks can then help in understanding the re-
 379 sponse of the large-scale circulation to our surface evaporative forcing. First, by verti-
 380 cally integrating the moist static energy budget, one can relate large-scale vertical winds
 381 to the local input of energy through the top and bottom boundaries of the atmosphere,

$$-\omega_1 M = NEI \quad (3)$$

382 where ω_1 is the vertical motion at a characteristic level, and M is a coefficient known
 383 as the gross moist stability (Neelin & Held, 1987; Sobel et al., 2007; Raymond et al., 2015).
 384 Being a measure of the vertical energy stratification of the atmosphere, M is typically
 385 positive in the time mean in deep-convecting regions, expressing the fact that time-mean
 386 ascent typically exports energy from the column in many tropical regions. Thus, the in-
 387 crease in NEI due to reduced low cloud cover is accompanied by enhanced large-scale
 388 ascent in many tropical land regions (Figs. 12a, S4). While this anomalous ascent is bottom-
 389 heavy, it extends the full depth of the troposphere. There is an increase in precipitation
 390 over many of the regions in which there is an increase in NEI (Fig. 12b). Outside of the
 391 tropics, precipitation over land decreases (Fig. 12b). Though there is increased time-mean
 392 upwards motion and increased NEI (Fig. S5), this does not lead to more precipitation,
 393 consistent with the fact that precipitation there is generated primarily by the moisture
 394 converged by transient motions, rather than time-mean flow.

395 In an alternate framework, one can consider the energy content of air (MSE) in-
 396 stead of the source of atmospheric energy (the NEI); precipitating ascent in the low-latitude
 397 atmosphere generally lies near the maximum in surface air MSE and increases in inten-
 398 sity with horizontal gradients in that MSE. This is expected when surface air MSE is

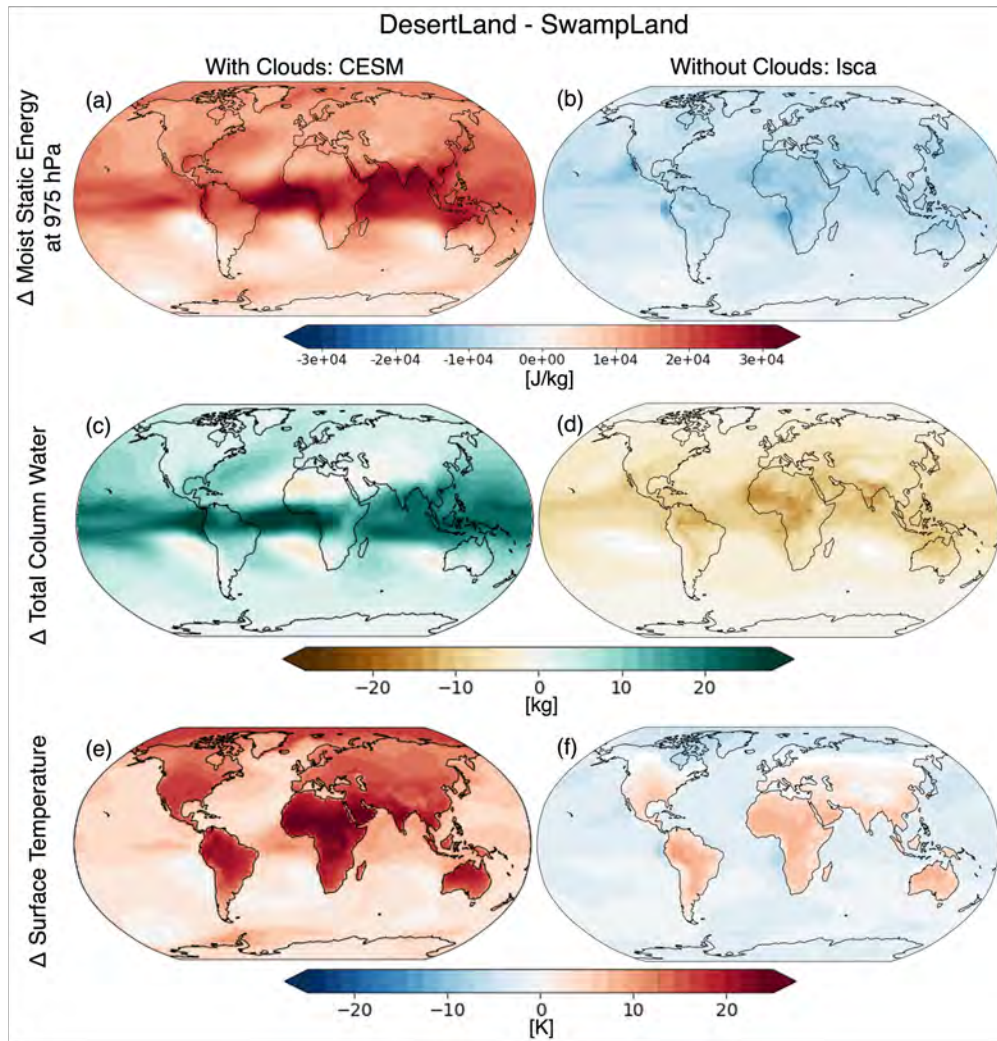


Figure 11. Annual mean change (DesertLand - SwampLand) in near-surface (975 hPa) moist static energy (top), total column water vapor (middle), and surface temperature (bottom) for CESM (left) and Isca (right) simulations. Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.

399 convectively coupled to free-tropospheric temperatures (Emanuel et al., 1994; Neelin, 2007;
 400 Privé & Plumb, 2007; Shekhar & Boos, 2016), and has been shown to describe the ob-
 401 served tropical climatology and interannual variability (Nie et al., 2010; Hurley & Boos,
 402 2013). Suppressed surface evaporation increases MSE (because of the atmospheric feed-
 403 backs discussed above), particularly over tropical regions which thus enhances the mag-
 404 nitude of the meridional gradient in MSE, and in turn the large-scale tropical overturn-
 405 ing circulation (Fig. 11a, 12). Decreases in moisture in the lower atmosphere over most
 406 land areas (driven by reduced surface evaporation) would lead to a reduction in MSE,
 407 but this is more than compensated for by increases in temperatures, which lead to an
 408 overall increase in MSE at 975 hPa over land areas (Fig. S6); changes in geopotential heights
 409 are negligible.

Neither of these frameworks provides a closure for the amount of water vapor in the atmosphere, but in Earth’s low-latitude atmosphere, where most water vapor lies and where horizontal temperature gradients are weak, MSE generally scales like precipitable water (Charney, 1963; Sobel et al., 2001) given the general availability of the ocean water supply in the modern continental configuration. In other words, large-scale ascent will advect water vapor upward, humidifying the column over land regions despite suppressed terrestrial evaporation. The MSE increase aloft over land in DesertLand indeed consists of increases in both temperatures and specific humidities (specific humidity decreases near the surface, but increases aloft), though the latter increase slowly enough that RH over land decreases (Fig. 5e); the increases in relative humidity occurring near the poles at high altitudes occur in the stratosphere and are not considered in this study (indeed, changes in T and Q above the tropopause are masked out of all calculations involving the radiative kernel, as in Pendergrass et al. (2018); Shell et al. (2008); Laguë, Swann, & Boos (2021)). While the warming in our simulations is driven by changes in land evaporation, for CO₂-driven warming, Byrne & O’Gorman (2016) also find near-surface continental relative humidity decreases. In simulations of future climate, we would expect both the radiative effects of CO₂ as well as changes in terrestrial evaporation to drive changes in atmospheric moisture. The only terrestrial regions where total column water vapor decreases is in the dry subtropics (e.g. the Sahara) where subsidence increases (Figs. S4, S5).

The discussion above differs greatly from arguments in which continental water vapor is treated as being set by transport from ocean regions. Anomalous near-surface winds do bring moist ocean air onto tropical land in DesertLand vs. SwampLand, e.g. in tropical South America, Africa, and Asia (Fig. 12b). However, such onshore winds need not produce ascent that spans most of the depth of the troposphere. Strong land-sea temperature gradients, such as those between a desert and ocean on Earth, often produce shallow, non-precipitating circulations; we will show below that, without cloud-radiative effects, suppressed land evaporation indeed leads to enhanced onshore flow that is shallow, non-precipitating, and that does not result in enhanced total column water vapor.

3.2.3 *The response without cloud-radiative effects*

When similar “desert” and “swamp” simulations are repeated with a cloud-free idealized global circulation model (Isca), surface temperatures still increase in response to suppressed terrestrial evaporation (Fig. 11f). However, both near-surface MSE (Fig. 11b) and atmospheric water vapor over continental regions (Fig. 11d) decrease. Unlike the CESM simulations, where the reduction in low cloud cover leads to an increase in energy into the Earth system at the top of the atmosphere, the Isca simulations do not have this additional energy source to the system as there are no clouds.

Suppressing terrestrial evaporation leads to a decrease in the atmospheric NEI and the net top of atmosphere energy budget over continents in the Isca simulations (Fig. 10). While the Isca simulations still show some anomalous upward motion over the continents in the lower atmosphere, and near-surface onshore flow from the oceans to the land in many regions (Figs. S7 & S8), these are shallow, non-precipitating anomalous circulations that contrast strongly with the deep, precipitating anomalous flow in CESM (c.f. S4 & S7). Similarly, the warming induced by suppressed terrestrial evaporation in Isca is restricted to the near-surface atmosphere over land regions, while CESM warms throughout the column (c.f. Fig. 5 & S9). While there are many differences between CESM and Isca, a key difference in the response to altered terrestrial evaporation is the response of cloud cover, which CESM includes and Isca does not. This highlights the importance of understanding cloud feedbacks—already a large source of climate uncertainty (Zelinka et al., 2017)—for determining how terrestrial evaporation changes alter the climate system.

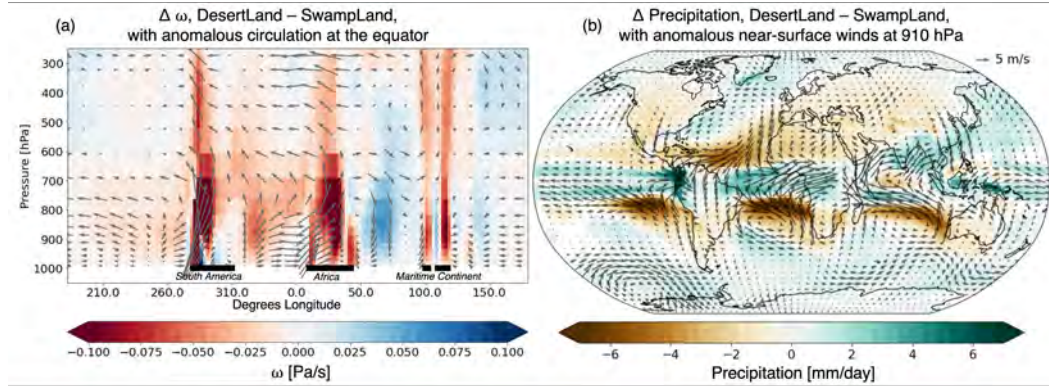


Figure 12. (a) Vertical cross section (longitude vs. pressure) at the equator for the change in vertical motion ω [Pa/s] for DesertLand – SwampLand is shown in colors, with red indicating negative ω , i.e. positive vertical motion. Vectors show the direction of anomalous motion: u winds in [m/s] in the x-direction, and $\omega \times -150$ in [Pa/s] in the y direction. ω is multiplied by -1 so arrows point in the direction of motion, and by 150 as an arbitrary scaling such that the vertical component is of comparable magnitude to the horizontal component, for ease of visualization. Horizontal black lines at the surface indicate land masses. (b) Change in precipitation (shading) for DesertLand – SwampLand, with vectors showing anomalous $\langle u, v \rangle$ winds [m/s] at 910 hPa. Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.

461 While we do not have a CESM simulation without radiatively interactive clouds,
 462 we can crudely compare the radiative fluxes in the Isca simulations to the top of atmo-
 463 sphere radiative fluxes calculated using "clear-sky" conditions in CESM (Fig. 11b/e).
 464 The clear-sky fluxes are calculated at each time step ignoring the radiative effects of clouds;
 465 however the model is integrated forwards using the full sky (including the radiative ef-
 466 fects of clouds) fluxes—that is, the temperature, moisture, dynamics, etc. of CESM are
 467 consistent with the full-sky radiative fluxes, not the clear-sky radiative fluxes. Both the
 468 clear-sky CESM simulations and the Isca simulations show a decrease in the TOA en-
 469 ergy balance over land regions when terrestrial evaporation is suppressed, reflecting the
 470 fact that the increase in energy into the Earth system in the CESM simulations are due
 471 to changes in low cloud cover over land.

472 While both CESM and Isca allow us to test the response of the climate system to
 473 suppressed terrestrial evaporation, we note that there are limitations in comparing re-
 474 sults between these two models, as they do not provide an apples-to-apples comparison.
 475 In future work, a perhaps more robust comparison would be to compare the full CESM-
 476 SLIM model to a version of CESM where the clouds have been modified to be transpar-
 477 ent to radiation; however, this model would produce unrealistically high TOA albedos,
 478 thus an increase in the surface albedo (such as is done in Isca) would be necessary to repli-
 479 cate a similar climate to the modern Earth. Such a modified model configuration would
 480 need to be evaluated and benchmarked before using it for this kind of experiment. Simi-
 481 larly, a modified version of Isca that included radiatively interactive clouds could also
 482 be leveraged to understand interactions between terrestrial evaporation, cloud cover, and
 483 the global energy budget. In this study however, we present the results of both the CESM
 484 and Isca simulations as complimentary; both indicate a strong control of terrestrial evap-
 485 oration on surface climate, and the model that allows for interactive cloud changes (CESM)
 486 suggests further investigation of the coupling of terrestrial evaporation and cloud cover
 487 is of merit. Additionally, we note that all the simulations shown in this study are highly

488 idealized; they are useful for improving our mechanistic understanding of interactions
489 between terrestrial processes and global climate, and can be used to inform more real-
490 istic studies of the effects of changes in terrestrial evapotranspiration driven by vegeta-
491 tion change, land use, agriculture, climate change, etc. on global climate.

492 4 Conclusions

493 In a global model with realistic continental geometry, reducing terrestrial evapo-
494 ration increases the total amount of atmospheric water vapor over most land and ocean
495 regions. The residence time of water vapor in the atmosphere increases by roughly 50%
496 from the simulation with fully saturated land to the simulation with desert land. Sup-
497 pressing land evaporation has a direct warming effect on the land surface by reducing
498 latent cooling of the surface, but also drives atmospheric feedbacks including reductions
499 in terrestrial cloud cover. The anomalous surface energy fluxes driven by atmospheric
500 cloud, water vapor, and temperature feedbacks are larger than the initial change in la-
501 tent heat flux driven directly by suppressed terrestrial evaporation. The cloud feedback
502 is critical for increasing near-surface moist static energy and generating anomalous at-
503 mospheric circulations throughout the depth of the troposphere. Simulations conducted
504 in a cloud-free model still show surface warming with suppressed terrestrial evaporation,
505 but also show a decrease, rather than an increase, in near surface MSE. Anomalous at-
506 mospheric circulations over the continents in cloud-free simulations are much shallower,
507 and the atmosphere shows reduced atmospheric water vapor with suppressed terrestrial
508 evaporation. This extreme experiment raises the question of how real-world changes to
509 the land surface (e.g. land use, agriculture) may be contributing to climate change by
510 altering atmospheric water vapor and cloud cover, and how terrestrial evaporation mod-
511 ulates climate on other planets or in past continental configurations of Earth's history.

512 Data Availability Statement

513 Source code for the models used are publicly available on github at [https://github](https://github.com/marysa/SimpleLand/releases/tag/slim0.1.006_release-cesm2.1.4)
514 [.com/marysa/SimpleLand/releases/tag/slim0.1.006_release-cesm2.1.4](https://github.com/marysa/SimpleLand/releases/tag/slim0.1.006_release-cesm2.1.4) for SLIM,
515 <https://github.com/ESCOMP/CESM/releases/tag/release-cesm2.1.1> for CESM, and
516 <https://github.com/ExeClim/Isca/releases/tag/v1.0> for Isca. The specific versions
517 used in this study are archived on zenodo in the following locations: SLIM, at [https://](https://zenodo.org/badge/latestdoi/509227365)
518 zenodo.org/badge/latestdoi/509227365 (DOI: 10.5281/zenodo.7352717); CESM, at
519 <https://doi.org/10.5281/zenodo.3895315> (DOI: 10.5281/zenodo.3895315); and Isca,
520 at <https://github.com/ExeClim/Isca/releases/tag/v1.0>, respectively. The mod-
521 ified version of Isca used in this study is available via zenodo at [https://doi.org/10](https://doi.org/10.5281/zenodo.6800218)
522 [.5281/zenodo.6800218](https://doi.org/10.5281/zenodo.6800218). Output from the simulations, analysis scripts, and model setup
523 files used in this study are available on Zenodo at [https://doi.org/10.5281/zenodo](https://doi.org/10.5281/zenodo.7909213)
524 [.7909213](https://doi.org/10.5281/zenodo.7909213)).

525 Author Contributions Statement

526 MML developed the idea for the study, modified the models, performed the sim-
527 ulations, and conducted the analysis. All authors discussed the experimental design and
528 results. MML wrote the initial draft and all authors contributed to the manuscript.

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1 **Supporting Information for: Reduced terrestrial**
2 **evaporation increases atmospheric water vapor by**
3 **generating cloud feedbacks**

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10 **Supplemental Figures**

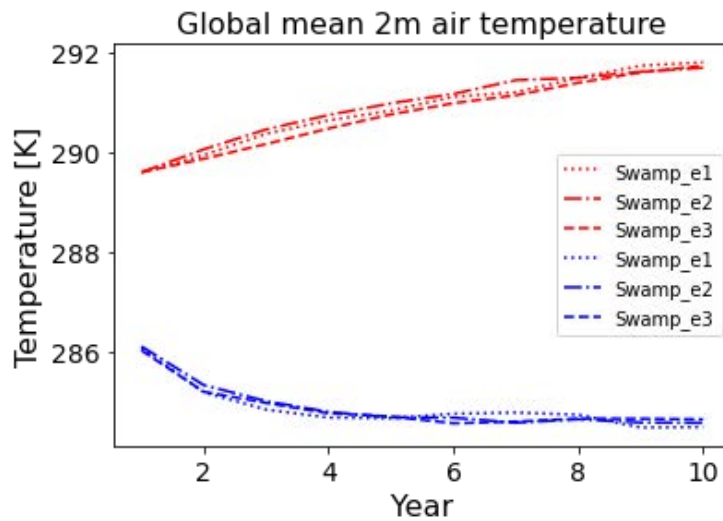


Figure S1. Spin-up of global mean 2m air temperatures across 3 ensemble members each of the Swamp and Desert simulations (swamp_e1-swamp_e3 and desert_e1-desert_e3, respectively), showing that the spread between ensemble members is small compared to inter-annual variability, and the difference between Swamp or Desert is much larger than the difference between ensemble members within either the Swamp or Desert simulation-type.

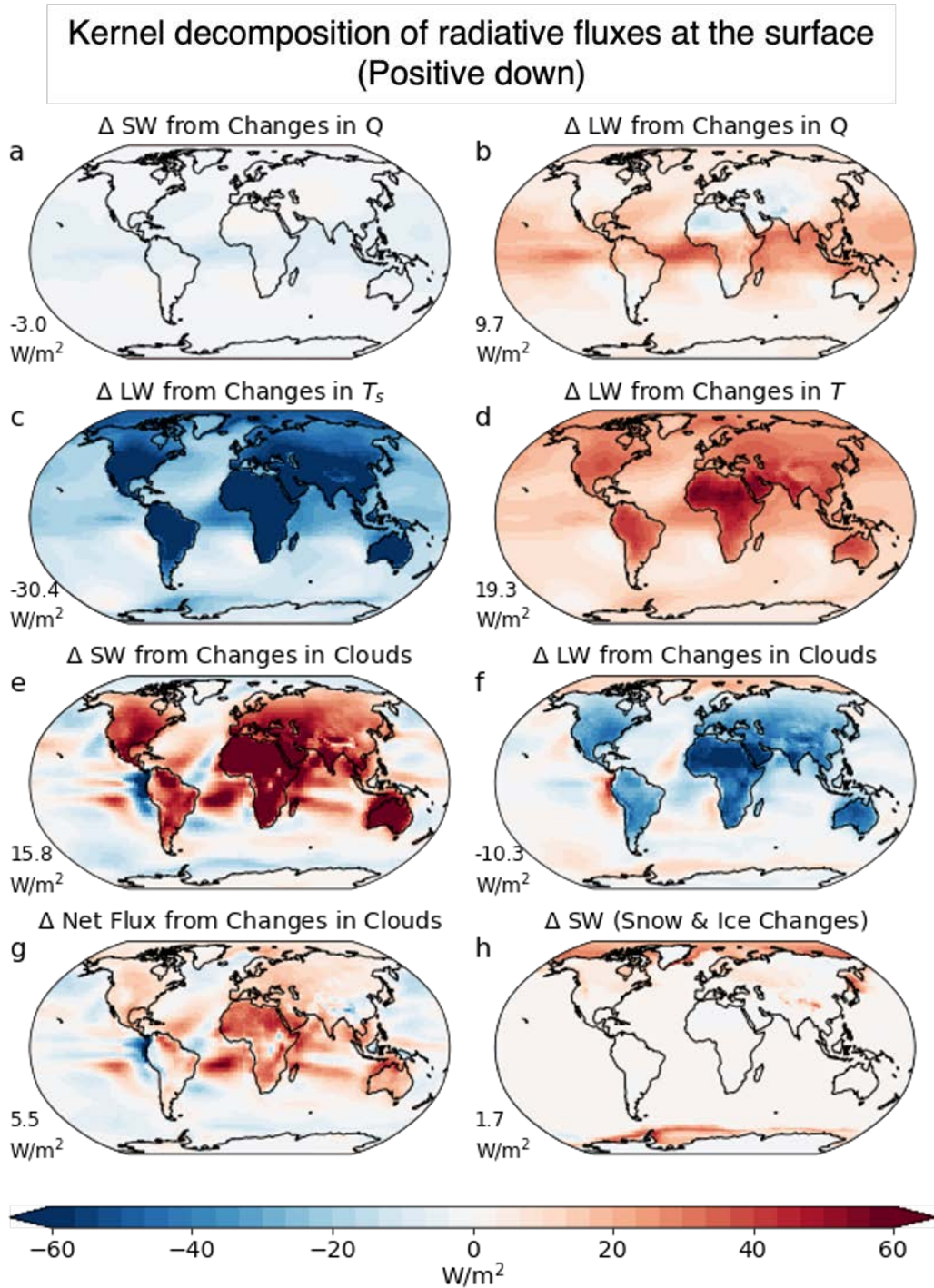


Figure S2. Decomposition of surface radiative fluxes using the radiative kernel for the DesertLand - SwampLand simulations. Fluxes are positive down into the surface, such that red regions indicate more energy into the surface attributable to that term in the DesertLand simulation than in the SwampLand simulation. The global mean value is noted to the lower left of each panel.

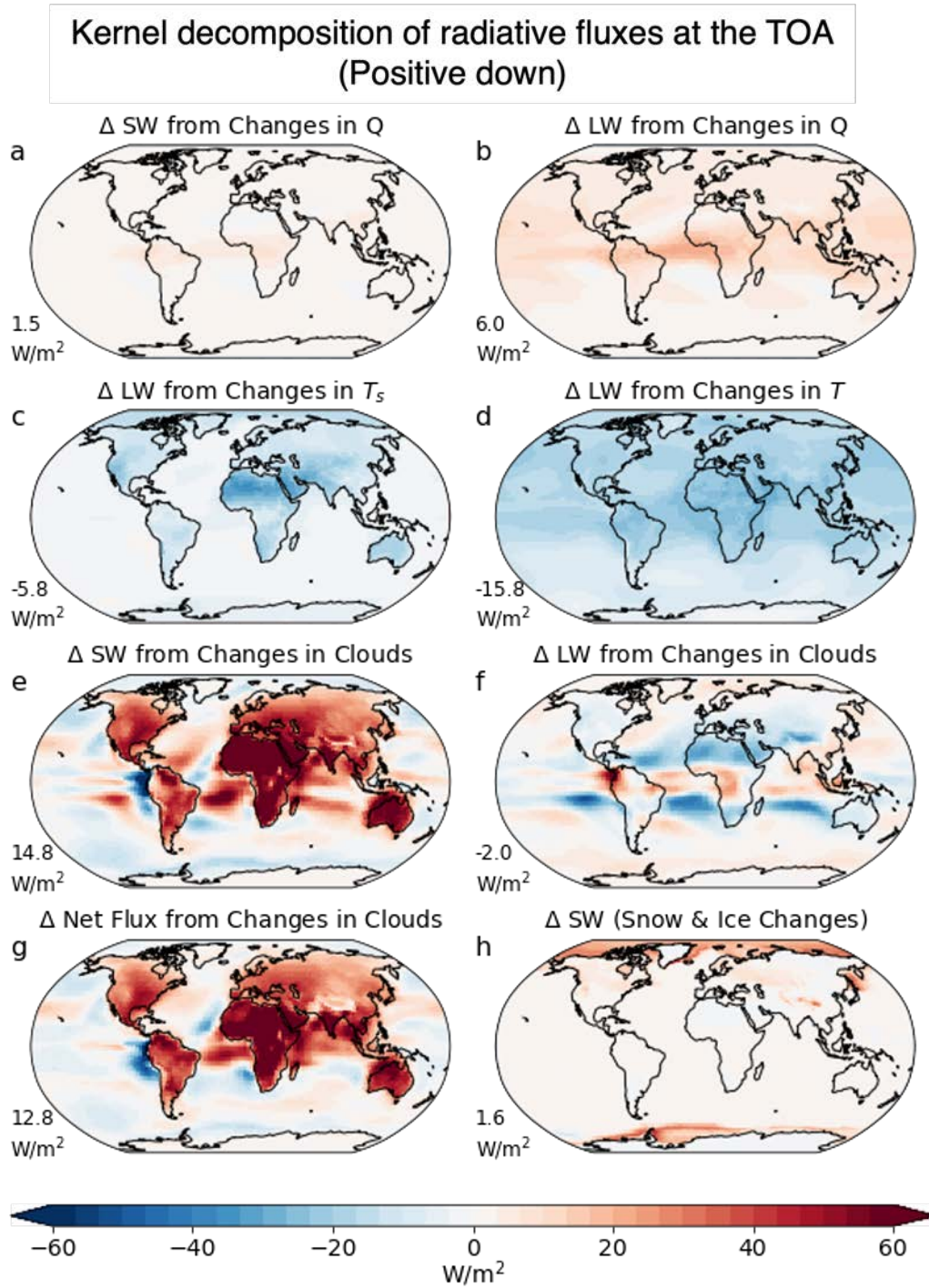


Figure S3. Decomposition of top of atmosphere radiative fluxes using the radiative kernel for the DesertLand - SwampLand simulations. Fluxes are positive down into the top of the atmosphere, such that red regions indicate more energy into the atmospheric column attributable to that term in the DesertLand simulation than in the SwampLand simulation. The global mean value is noted to the lower left of each panel.

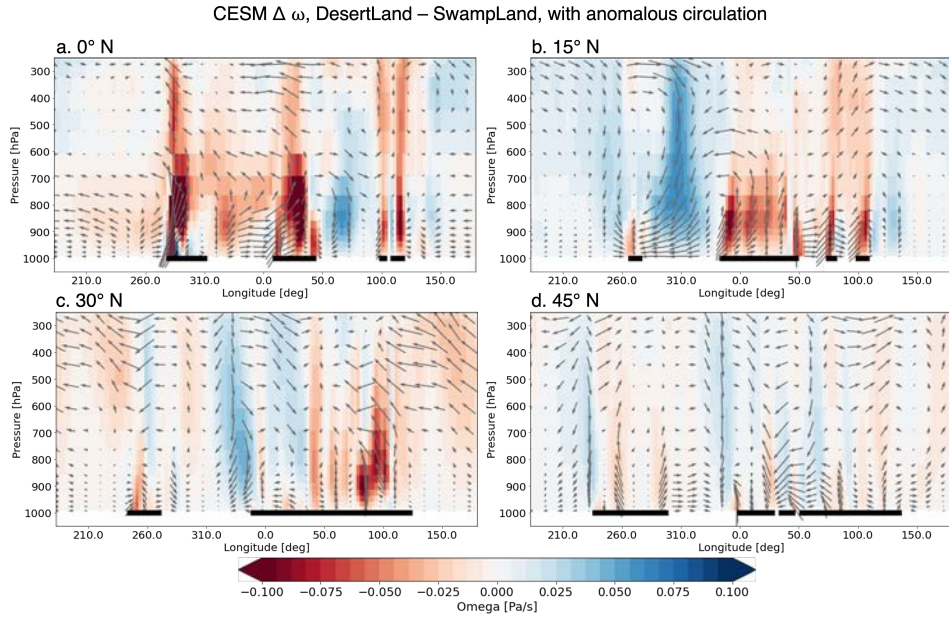


Figure S4. CESM vertical cross section (longitude vs. pressure) at various latitudes for the change in vertical motion ω [Pa/s] for DesertLand – SwampLand is shown in colors, with red indicating negative ω , i.e. positive vertical motion. Vectors show the direction of anomalous motion: u winds in [m/s] in the x-direction, and $\omega \times -150$ in [Pa/s] in the y direction. ω is multiplied by -1 so arrows point in the direction of motion, and by 50 so the vertical component is of comparable magnitude to the horizontal component, for ease of visualization. Horizontal black lines at the surface indicate land masses.

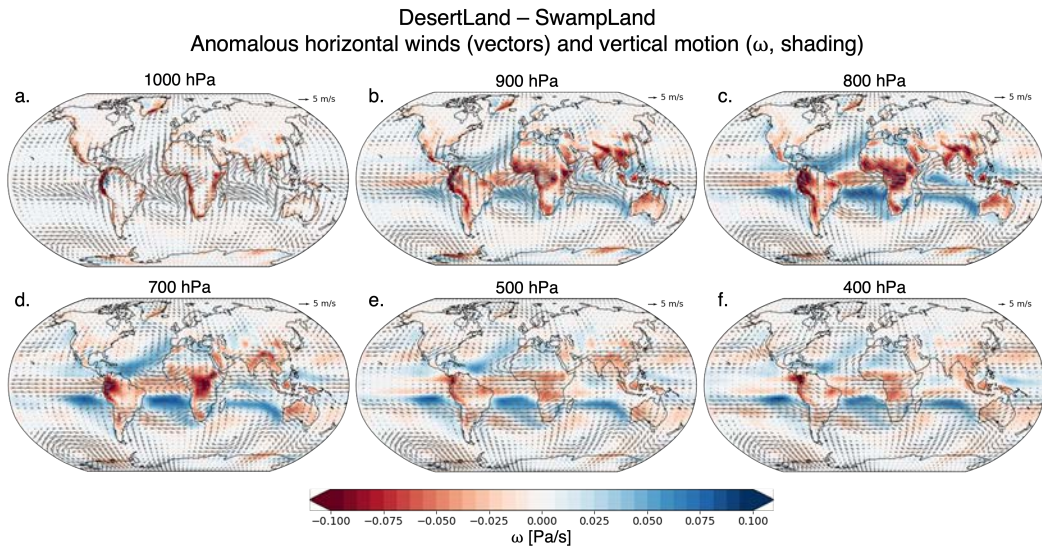


Figure S5. CESM change (DesertLand – SwampLand) in vertical motion ω [Pa/s] (shading, with red indicating upwards motion) and horizontal winds $\langle u, v \rangle$ [m/s] (vectors) at various levels of the atmosphere.

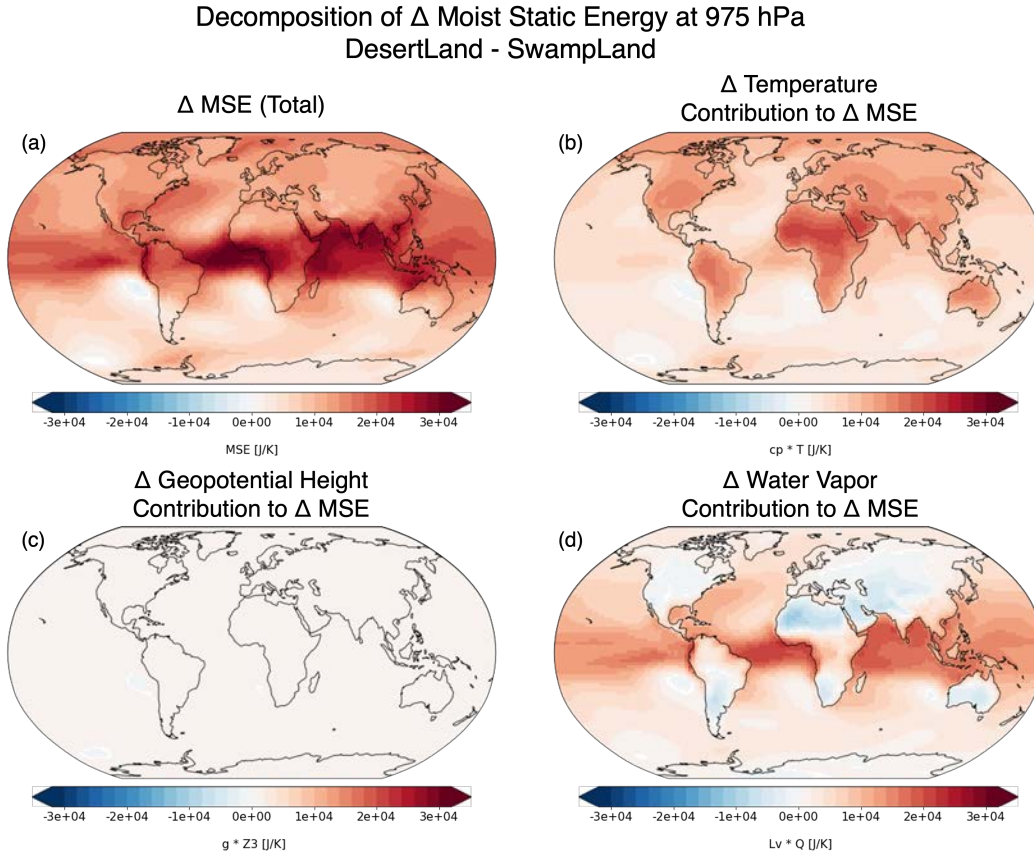


Figure S6. Decomposition of the change in moist static energy (MSE) into (a) the total and the components attributable to the change in (b) temperatures ($c_p T$), (c) geopotential heights (gZ), and (d) moisture ($L_v Q$) at 975 hPa, for the difference between the Desert and Swamp CESM simulations. Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.

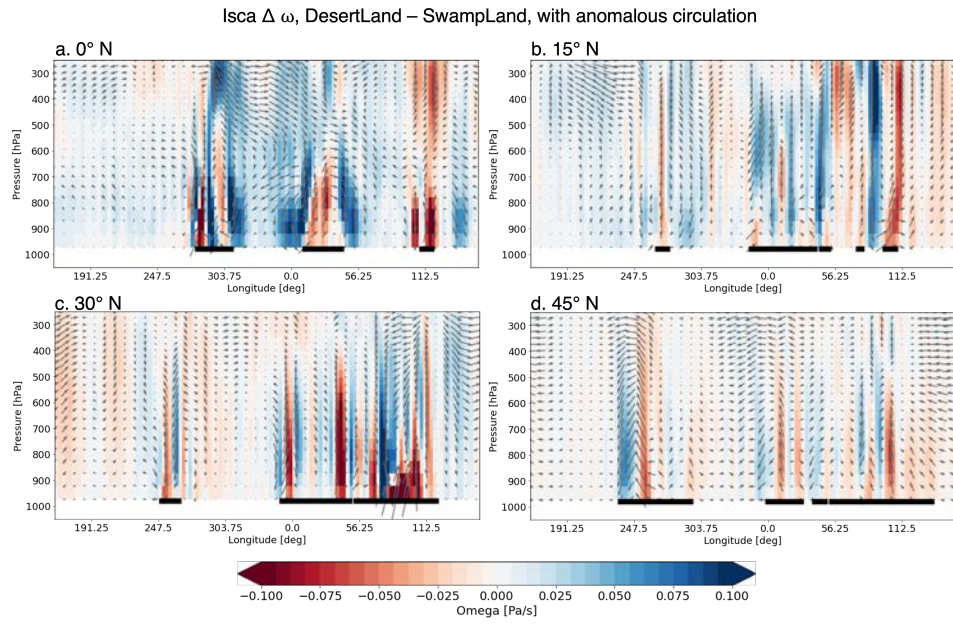


Figure S7. As in Fig. ??, but for Isca.

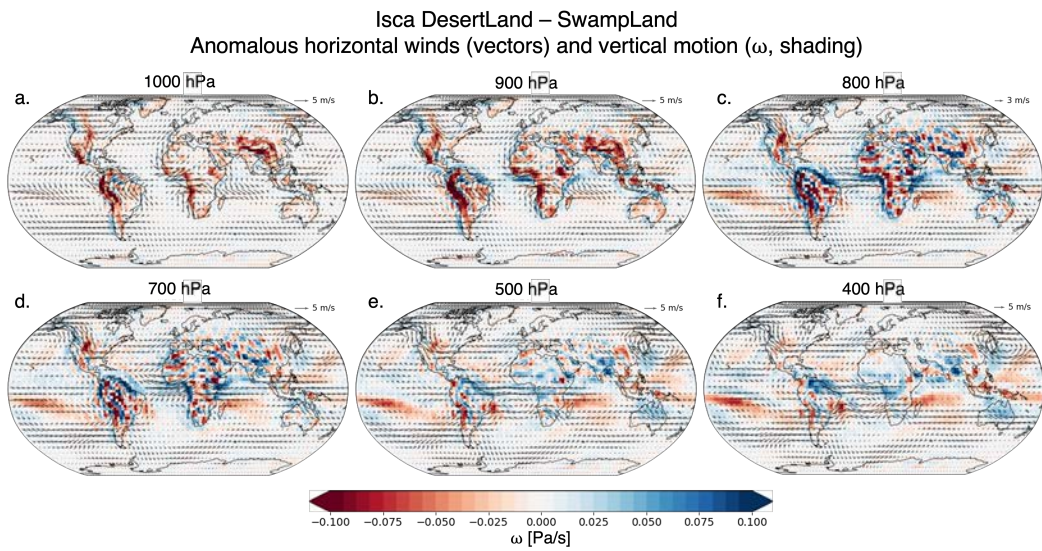


Figure S8. As in Fig. ??, but for Isca.

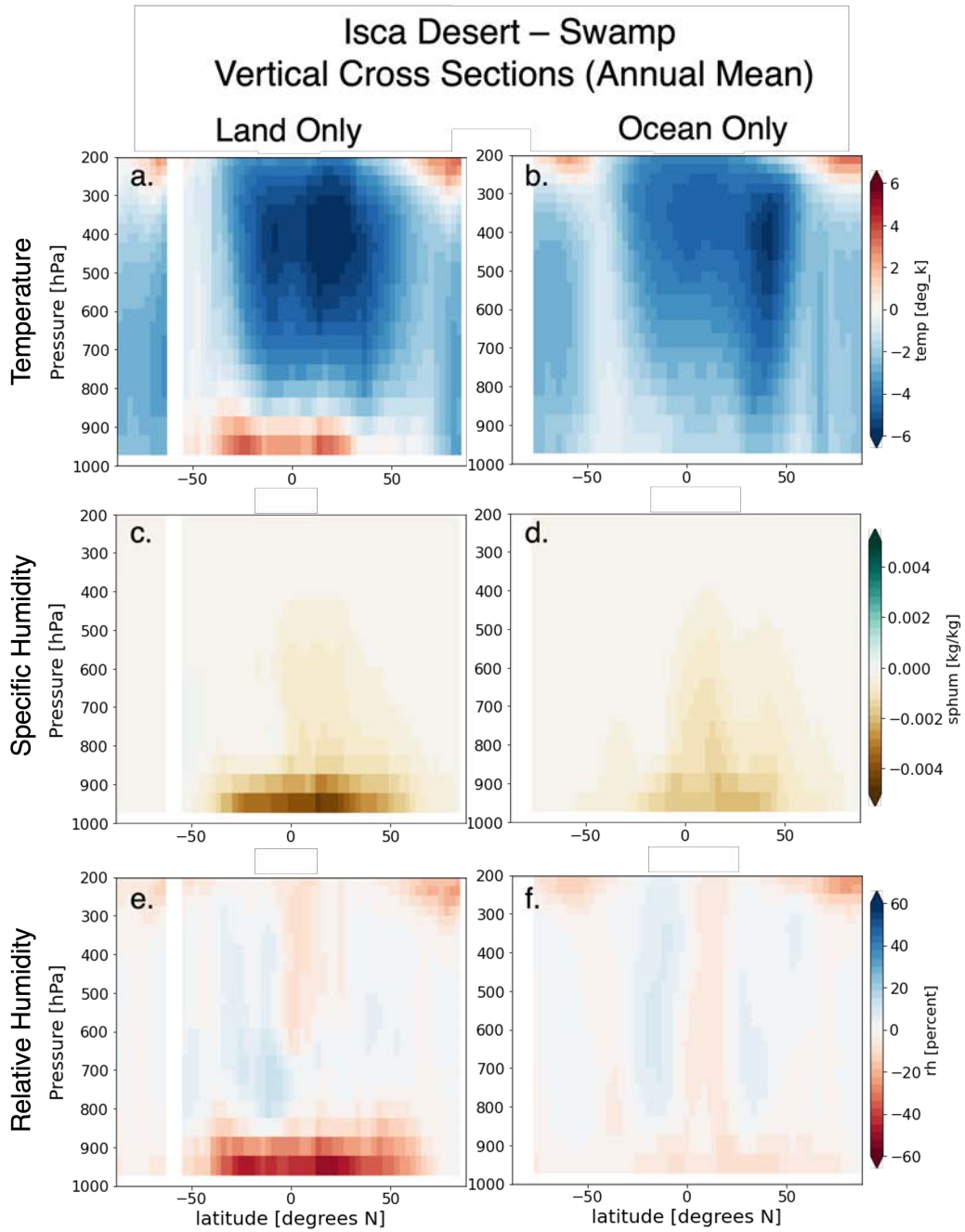


Figure S9. Change in the Isca simulations (DesertLand - SwampLand) in zonally averaged temperature (top), specific humidity (middle) and relative humidity (bottom) for land-only (left) and ocean-only (right).